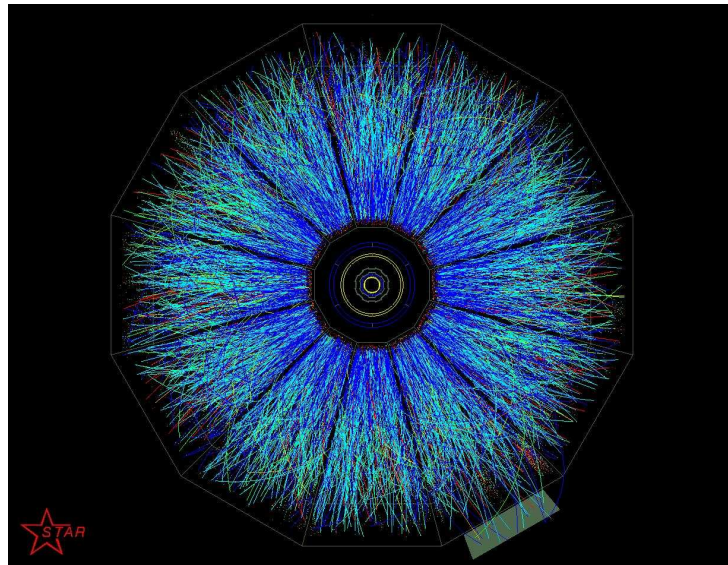
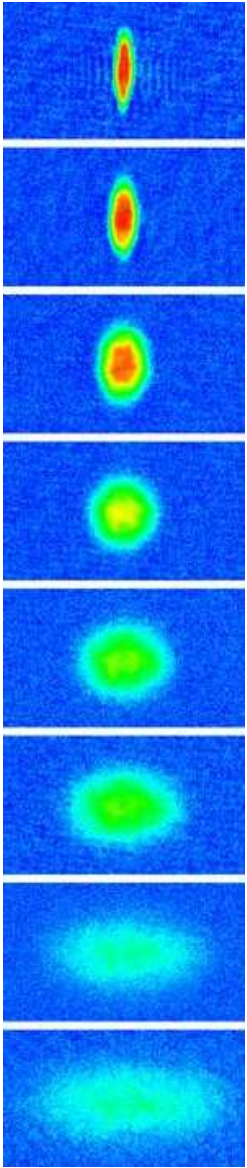


# Perfect Fluid Olympics

Thomas Schaefer

North Carolina State University

# Perfect Fluids: The contenders



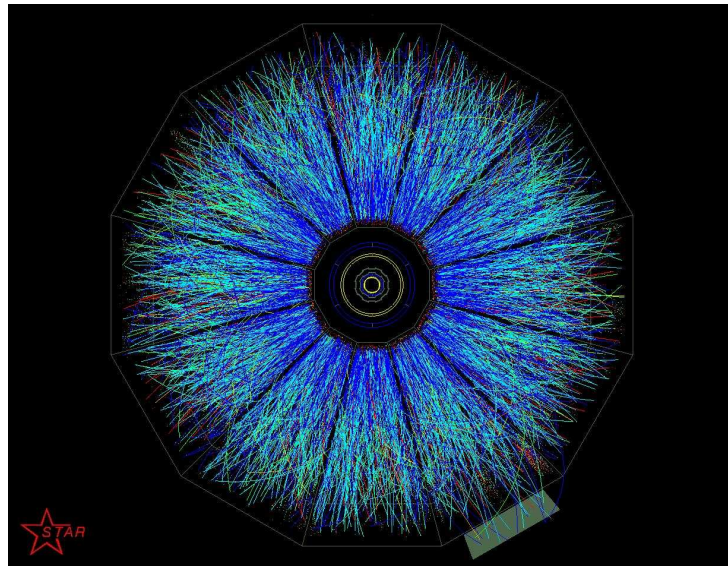
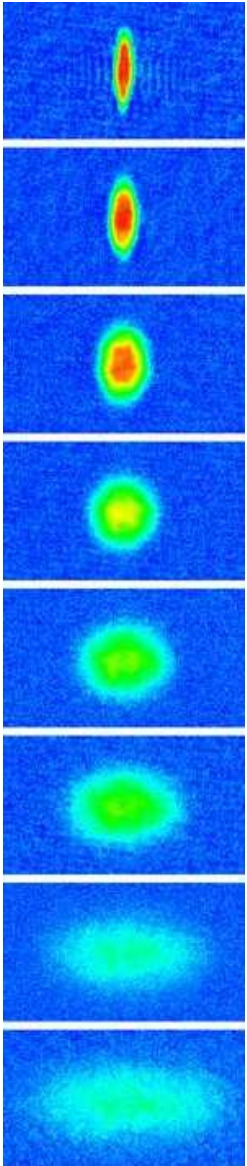
QGP ( $T=180$  MeV)

Trapped Atoms  
( $T=0.1$  neV)



Liquid Helium  
( $T=0.1$  meV)

# Perfect Fluids: The contenders



QGP  $\eta = 5 \cdot 10^{11} Pa \cdot s$

Trapped Atoms

$\eta = 1.7 \cdot 10^{-15} Pa \cdot s$



Liquid Helium

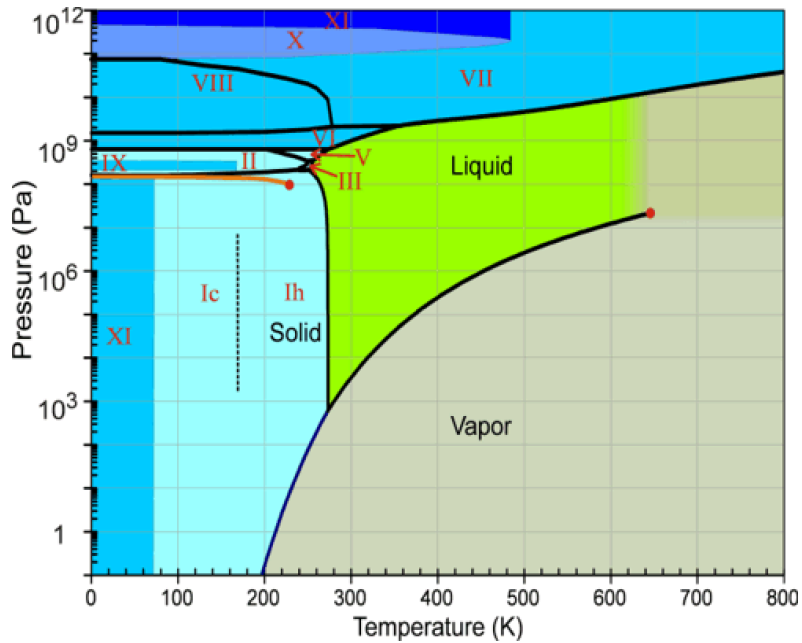
$\eta = 1.7 \cdot 10^{-6} Pa \cdot s$

Consider ratios

$\eta/s$

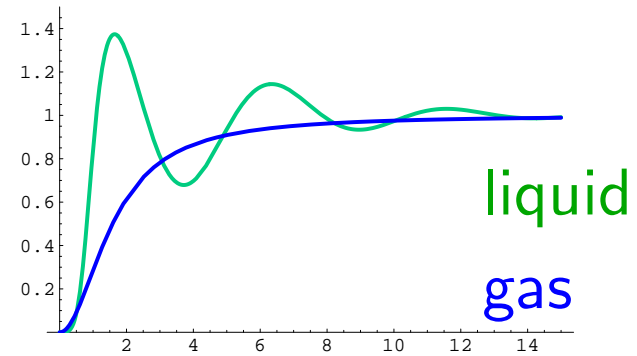
# Transitions without change of symmetry: Liquid-Gas

Phase diagram of water

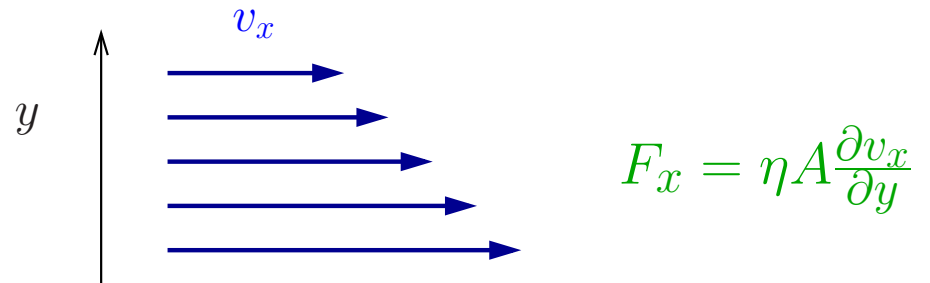


Characteristics of a liquid

Pair correlation function

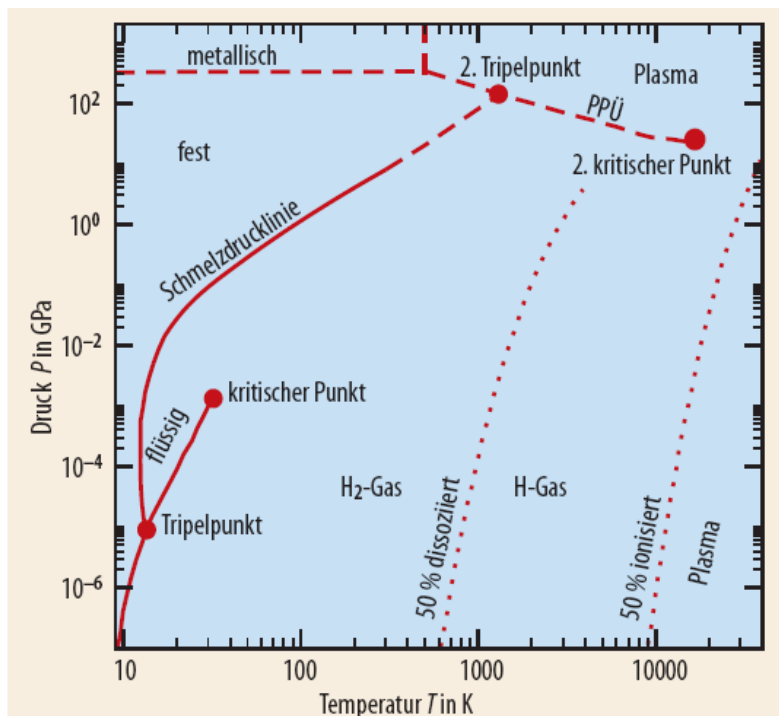


Good fluid: low viscosity



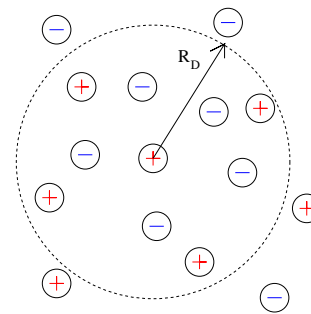
# Transitions without change of symmetry: Gas-Plasma

## Phase diagram of hydrogen



## Plasma Effects

### Debye screening



$$V(r) = -\frac{e}{r} e^{-m_D r}$$

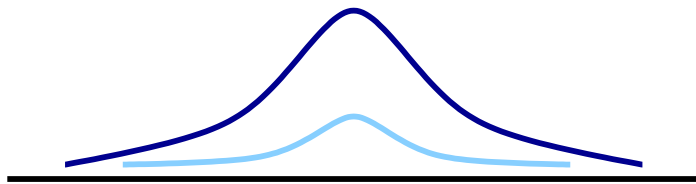
$$m_D^2 = \frac{4\pi e^2 n}{kT}$$

### Plasma oscillations

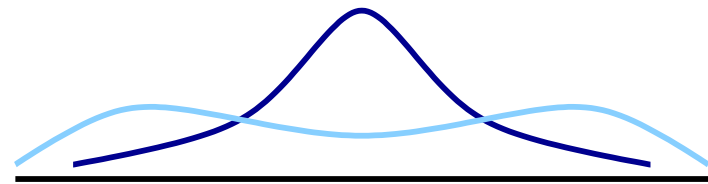
$$\omega_{pl} = \frac{4\pi e^2 n}{m}$$

# Fluids: Gases, Liquids, Plasmas, ...

Hydrodynamics: Long-wavelength, low-frequency dynamics of conserved or spontaneously broken symmetry variables.



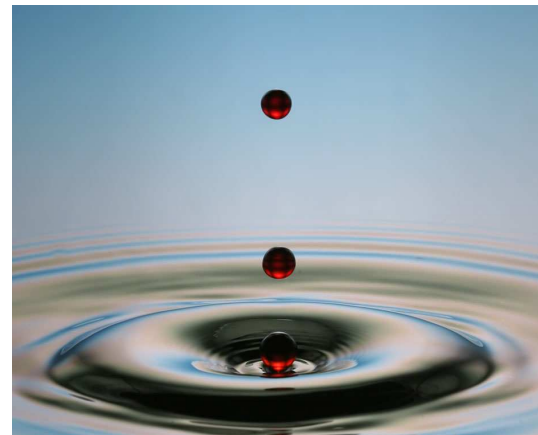
$$\tau \sim \tau_{micro}$$



$$\tau \sim \lambda^{-1}$$

Historically: Water

$$(\rho, \epsilon, \vec{\pi})$$



# Example: Simple Fluid

Continuity equation

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

Euler (Navier-Stokes) equation

$$\frac{\partial}{\partial t} (\rho v_i) + \frac{\partial}{\partial x_j} \Pi_{ij} = 0$$



Energy momentum tensor

$$\Pi_{ij} = P \delta_{ij} + \rho v_i v_j + \eta \left( \partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right) + \dots$$

reactive

dissipative

# Kinetic Theory

Quasi-Particles ( $\gamma \ll \omega$ ): introduce distribution function  $f_p(x, t)$

$$N = \int d^3p f_p \quad T_{ij} = \int d^3p \frac{p_i p_j}{E_p} f_p,$$

Boltzmann equation

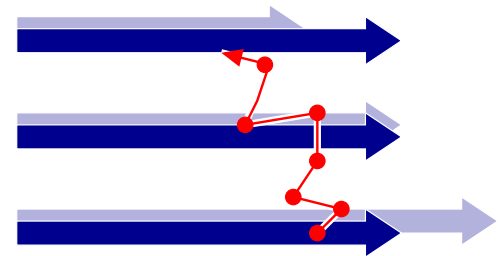
$$\frac{\partial f_p}{\partial t} + \vec{v} \cdot \vec{\nabla}_x f_p + \vec{F} \cdot \vec{\nabla}_p f_p = C[f_p]$$

Collision term  $C[f_p] = C_{gain} - C_{loss}$

Linearized theory (Chapman-Enskog):  $f_p = f_p^0 (1 + \chi_p/T)$

suitable for transport coefficients

shear viscosity  $\chi_p = g_p p_i p_j v_{ij}$





## Viscosity Bound: Rough Argument

Kinetic theory estimate of shear viscosity

$$\eta \sim \frac{1}{3} n \bar{p} l_{mfp} \quad (\text{Note : } l_{mfp} \sim 1/(n\sigma))$$

Normalize to density. Uncertainty relation implies

$$\frac{\eta}{n} \sim \frac{n \bar{p} l_{mfp}}{n} \geq \hbar$$

Also:  $s \sim k_B n$  and  $\eta/s \geq \hbar/k_B$

Validity of kinetic theory as  $\bar{p} l_{mfp} \sim \hbar$ ?

What is the “right” ratio,  $\eta/s$  or  $\eta/n$ ?

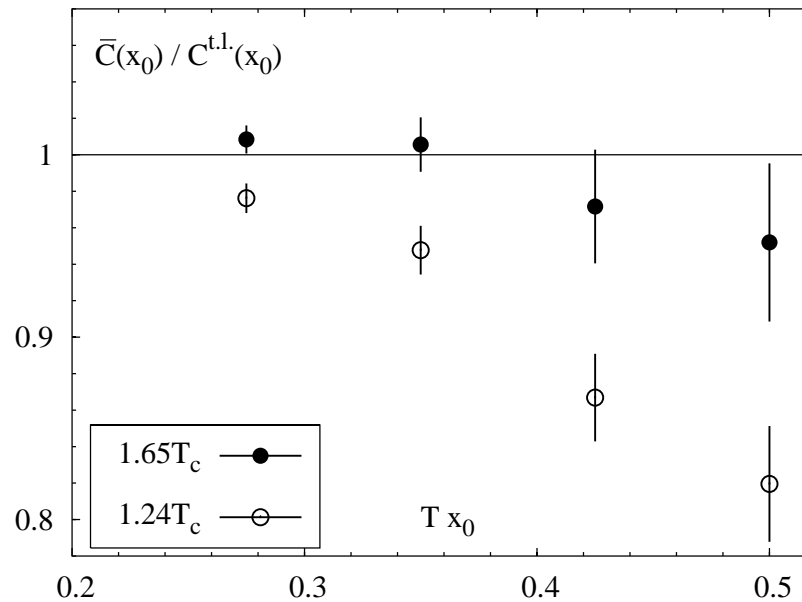
# Kubo Formula

Linear response theory provides relation to Green functions

$$\eta = - \lim_{\omega \rightarrow 0} \frac{1}{\omega} G_R(\omega, 0)$$

$$G_R(\omega, 0) = \int dt d^3x e^{i\omega t} \Theta(t) \langle [T_{xy}(t, x), T_{xy}(0, 0)] \rangle$$

Can be used in field theory (hard), lattice calculations (also hard)



## Dynamic Universality

Continuous phase transition: Dynamics of low energy modes universal

### Universality for transport coefficients

Universal theory: Hydro (diffusive modes), order parameters (time dependent LG), stochastic forces (Langevin)

$$\frac{\partial}{\partial t}(\rho v_i) = P_{ij}^\perp \left[ \eta_0 \nabla^2 \frac{\delta \mathcal{H}}{\delta(\rho v_j)} + w_0 (\nabla_j \phi) \frac{\delta \mathcal{H}}{\delta \phi} + \zeta_j \right]$$

Model H of Hohenberg and Halperin

$$\eta \sim \xi^{x_\eta} \quad (x_\eta \simeq 0.06) \qquad \zeta \sim \xi^{x_\zeta} \quad (x_\zeta \simeq 2.8)$$

# Gauge Theory at Strong Coupling: Holographic Duals

The AdS/CFT duality relates

large  $N_c$  (Conformal) gauge  
theory in 4 dimensions



string theory on 5 dimensional  
Anti-de Sitter space  $\times S_5$

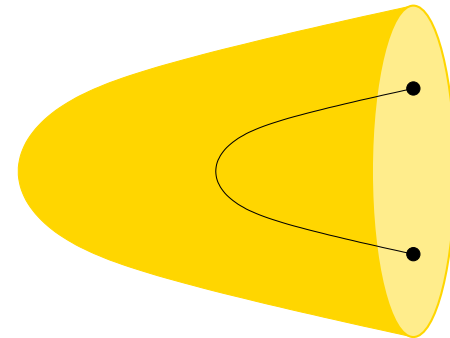
correlation fcts of gauge  
invariant operators



boundary correlation fcts  
of AdS fields

$$\langle \exp \int dx \phi_0 \mathcal{O} \rangle =$$

$$Z_{string}[\phi(\partial AdS) = \phi_0]$$



The correspondence is simplest at strong coupling  $g^2 N_c$

strongly coupled gauge theory  $\Leftrightarrow$

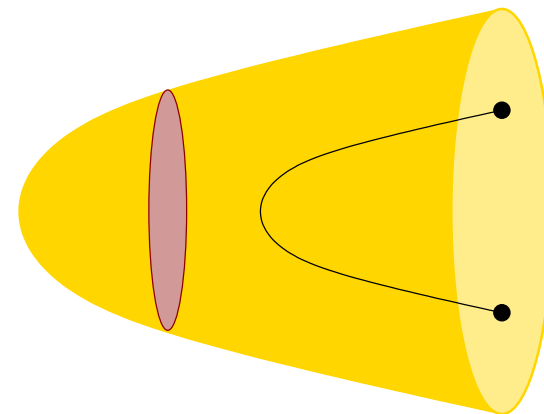
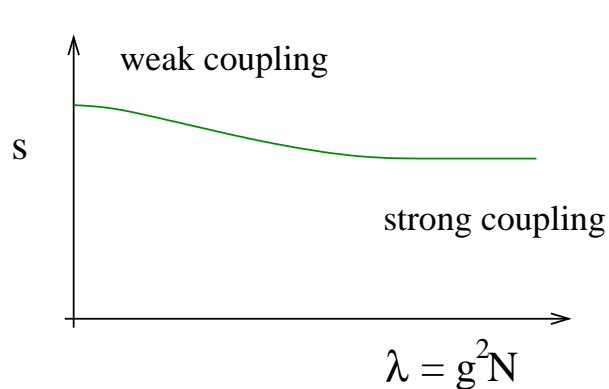
classical string theory

# Holographic Duals at Finite Temperature

Thermal (conformal) field theory  $\equiv AdS_5$  black hole

CFT temperature  $\Leftrightarrow$  Hawking temperature of black hole

CFT entropy  $\Leftrightarrow$  Hawking-Bekenstein entropy  
 $\sim$  area of event horizon



$$s(\lambda \rightarrow \infty) = \frac{\pi^2}{2} N_c^2 T^3 = \frac{3}{4} s(\lambda = 0)$$

# Holographic Duals: Transport Properties

Thermal (conformal) field theory  $\equiv AdS_5$  black hole

CFT entropy  $\Leftrightarrow$

Hawking-Bekenstein entropy

$\sim$  area of event horizon

shear viscosity  $\Leftrightarrow$

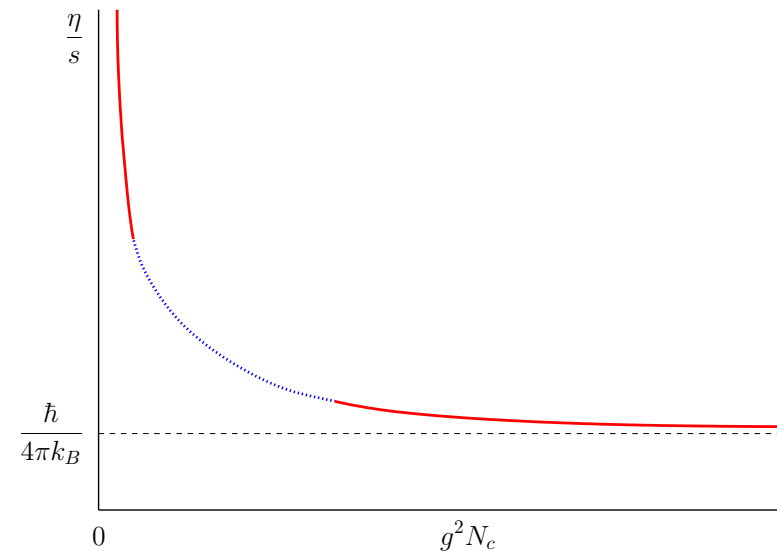
Graviton absorption cross section

$\sim$  area of event horizon

Strong coupling limit

$$\frac{\eta}{s} = \frac{\hbar}{4\pi k_B}$$

Son and Starinets



Strong coupling limit universal? Provides lower bound for all theories?

# Holographic Duals: Recent Work

Can we get  $AdS/NRCFT$  using an embedding of non-relativistic conformal group?

$$Schr(d) \subset SO(d+1, 2) \quad (= Iso(AdS_{d+2}))$$

Not quite, but we can use  $Schr(d) \subset SO(d+2, 2)$

$$\text{New } 2+1 \text{ NRCFTs with } \frac{\eta}{s} = \frac{1}{4\pi}$$

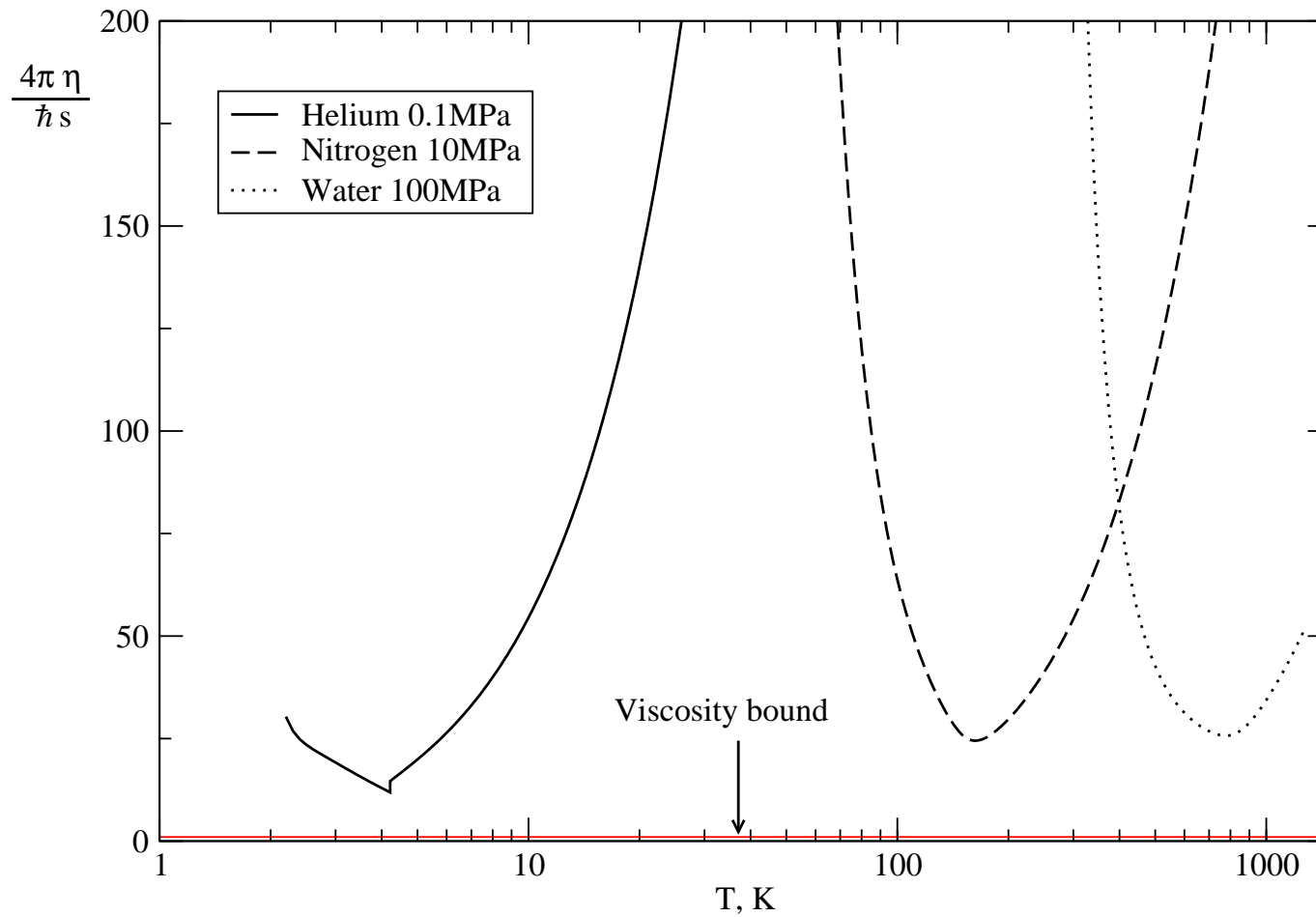
Son, McGreevy, Balasubramanian, . . .

Can we prove viscosity bound? Proven in certain classes of backgrounds, but some possible counterexamples exist

$$\frac{\eta}{s} = \frac{1}{4\pi} (1 - 4\lambda_{GB}) \quad \lambda_{GB} < \frac{9}{100}$$

Myers, Shenker, Kats, Petrov, . . .

# Viscosity Bound: Common Fluids





## Perfect Fluids: How to be a contender?

Bound is quantum mechanical

need quantum fluids

Bound is incompatible with weak coupling and kinetic theory

strong interactions, no quasi-particles

Model system has conformal invariance (essential?)

(Almost) scale invariant systems

# I. Fermi gas at unitarity

Non-relativistic fermions at low momentum

$$\mathcal{L}_{\text{eff}} = \psi^\dagger \left( i\partial_0 + \frac{\nabla^2}{2M} \right) \psi - \frac{C_0}{2} (\psi^\dagger \psi)^2$$

Unitary limit:  $a \rightarrow \infty, \sigma \rightarrow 4\pi/k^2$  ( $C_0 \rightarrow \infty$ )

This limit is smooth (HS-trafo,  $\Psi = (\psi_\uparrow, \psi_\downarrow^\dagger)$ )

$$\mathcal{L} = \Psi^\dagger \left[ i\partial_0 + \sigma_3 \frac{\vec{\nabla}^2}{2m} \right] \Psi + (\Psi^\dagger \sigma_+ \Psi \phi + h.c.) - \frac{1}{C_0} \phi^* \phi ,$$

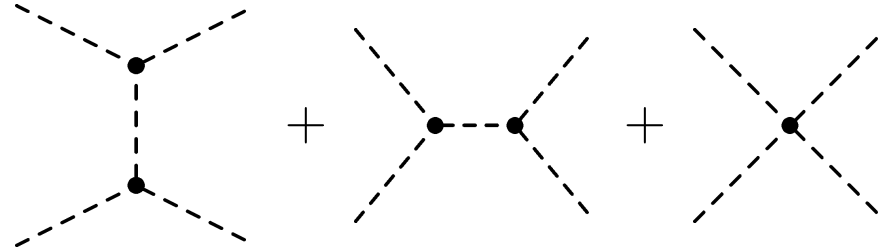
Low  $T$  ( $T < T_c \sim \mu$ ): Pairing and superfluidity

Low T: Phonons Goldstone boson  $\psi\psi = e^{2i\varphi} \langle \psi\psi \rangle$

$$\mathcal{L} = c_0 m^{3/2} \left( \mu - \dot{\varphi} - \frac{(\vec{\nabla}\varphi)^2}{2m} \right)^{5/2} + \dots$$

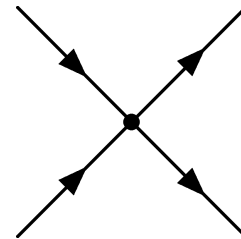
Viscosity dominated by  $\varphi + \varphi \rightarrow \varphi + \varphi$

$$\eta = A \frac{\xi^5}{c_s^3} \frac{T_F^8}{T^5}$$



High T: Atoms Cross section regularized by thermal momentum

$$\eta = \frac{15}{32\sqrt{2}} (mT)^{3/2}$$



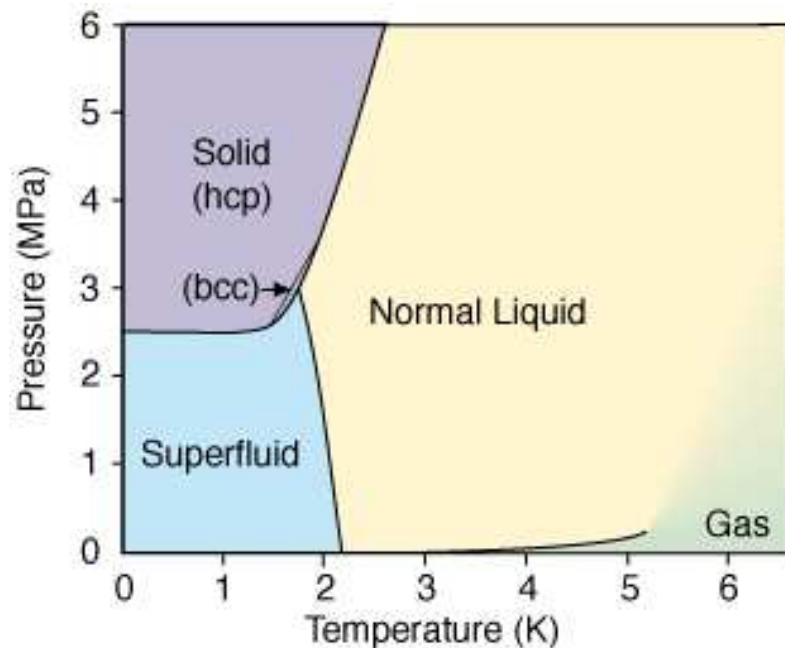
## II. Liquid Helium

Bosons, van der Waals + short range repulsion

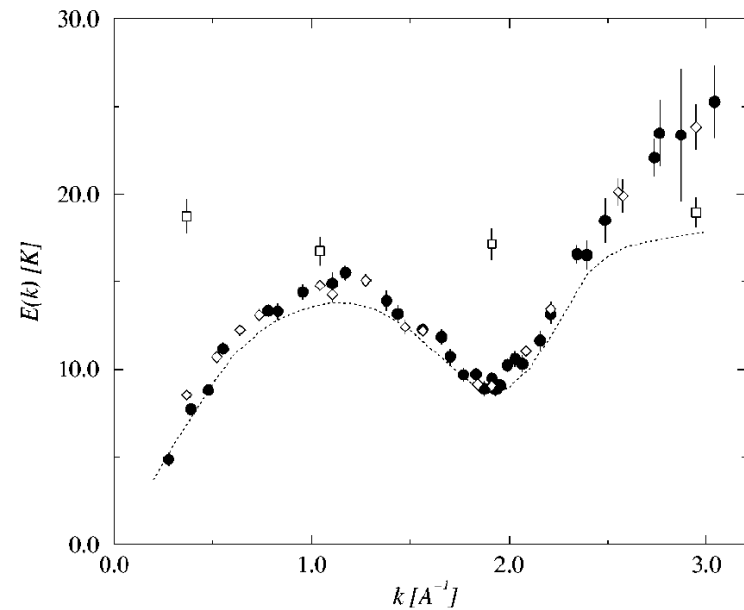
$$S = \int \Phi^\dagger \left( i\partial_0 + \frac{\nabla^2}{2M} \right) \Phi + \int \int (\Phi^\dagger \Phi) V(x - y) (\Phi^\dagger \Phi)$$

with  $V(x) = V_{sr}(x) - c_6/x^6$ . Note:  $a = 189a_0 \gg a_0$

### Phase Diagram



### Excitations



## Low T: Phonons and Rotons Effective lagrangian

$$\begin{aligned}\mathcal{L} = & \varphi^* (\partial_0^2 - v^2) \varphi + i\lambda \dot{\varphi} (\vec{\nabla} \varphi)^2 + \dots \\ & + \varphi_{R,v}^* (i\partial_0 - \Delta) \varphi_{R,v} + c_0 (\varphi_{R,v}^* \varphi_{R,v})^2 + \dots\end{aligned}$$

Shear viscosity

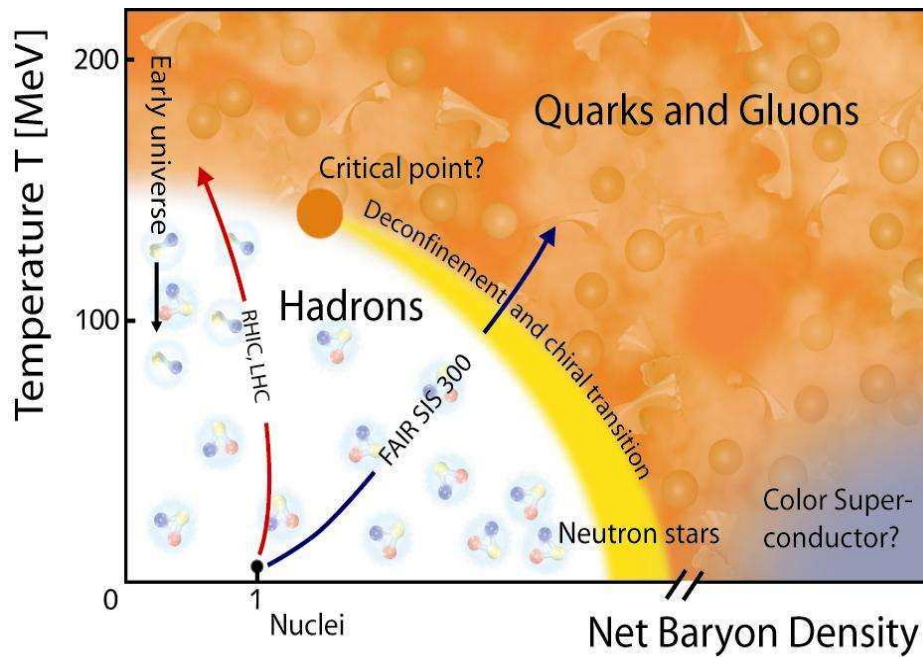
$$\eta = \eta_R + \frac{c_{Rph}}{\sqrt{T}} e^{\Delta/T} + \frac{c_{ph}}{T^5} + \dots$$

High T: Atoms Viscosity governed by hard core ( $V \sim 1/r^{12}$ )

$$\eta = \eta_0 (T/T_0)^{2/3}$$

# III. Quark Gluon Plasma

$$\mathcal{L} = \bar{q}_f (i\not{D} - m_f) q_f - \frac{1}{4} G_{\mu\nu}^a G_{\mu\nu}^a$$

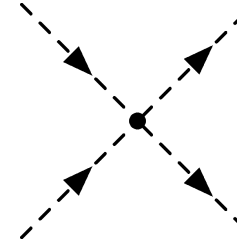


## Low T: Pions Chiral perturbation theory

$$\mathcal{L} = \frac{f_\pi^2}{4} \text{Tr}[\partial_\mu U \partial^\mu U^\dagger] + (B \text{Tr}[MU] + h.c.) + \dots$$

Viscosity dominated by  $\pi\pi$  scattering

$$\eta = A \frac{f_\pi^4}{T}$$



## High T: Quasi-Particles HTL theory (screening, damping, ...)

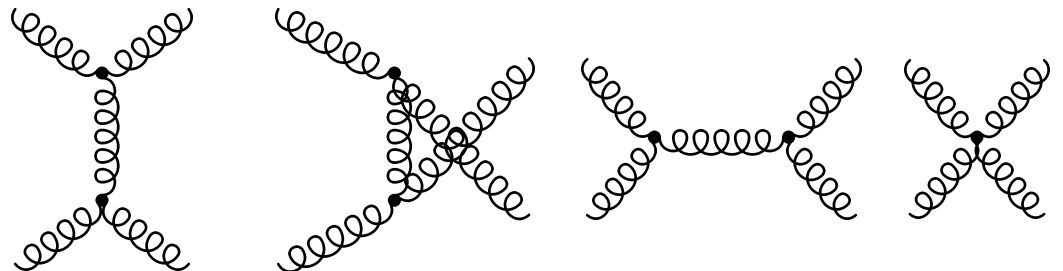
$$\mathcal{L}_{HTL} = \int d\Omega G_{\mu\alpha}^a \frac{v^\alpha v_\beta}{(v \cdot D)^2} G^{a,\mu\beta}$$

quasi-particle width

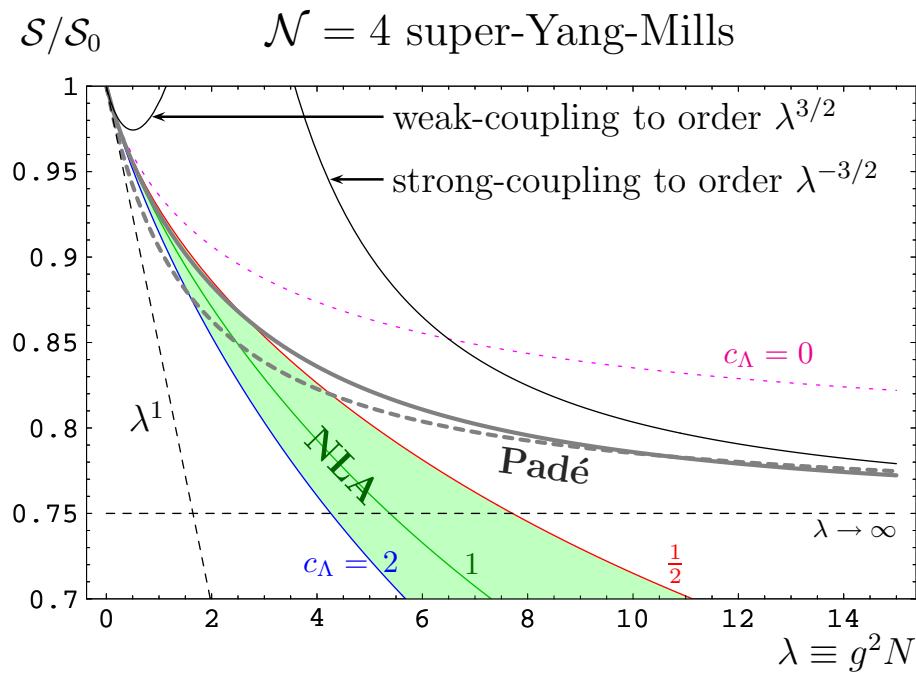
$$\gamma \sim g^2 T$$

Viscosity dominated by t-channel gluon exchange

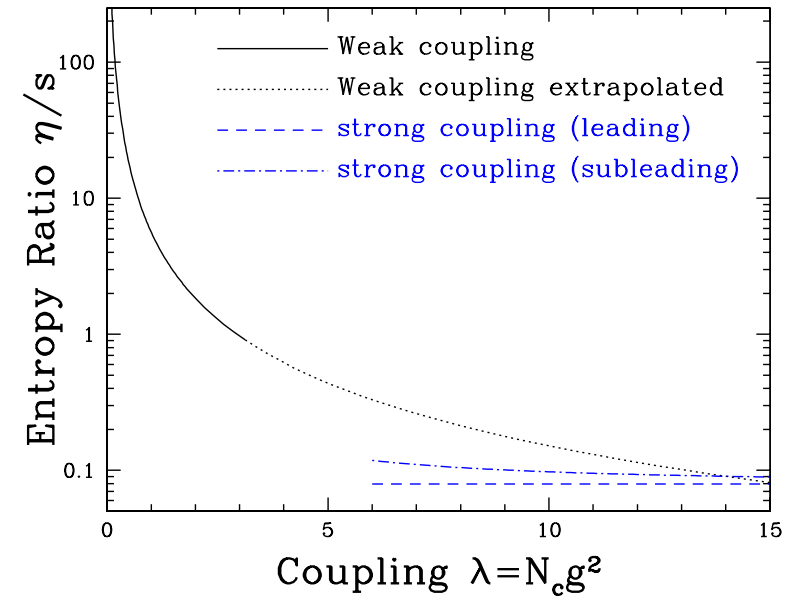
$$\eta = \frac{27.13 T^3}{g^4 \log(2.7/g)}$$



# What is the value of the coupling?



Blaizot, Iancu, Rebhan, Kraemmer, Rebhan (2006)



Huot, Jeon, Moore (2006)



# I. Experiment (Liquid Helium)

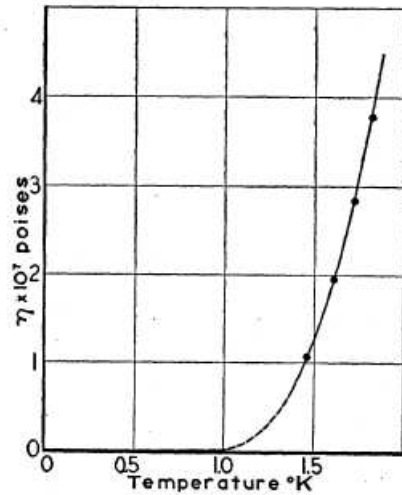


FIG. 1. The viscosity of liquid helium II measured by flow through a  $10^{-4}$  cm channel.

2, 53

LIQUID HELIUM II

23

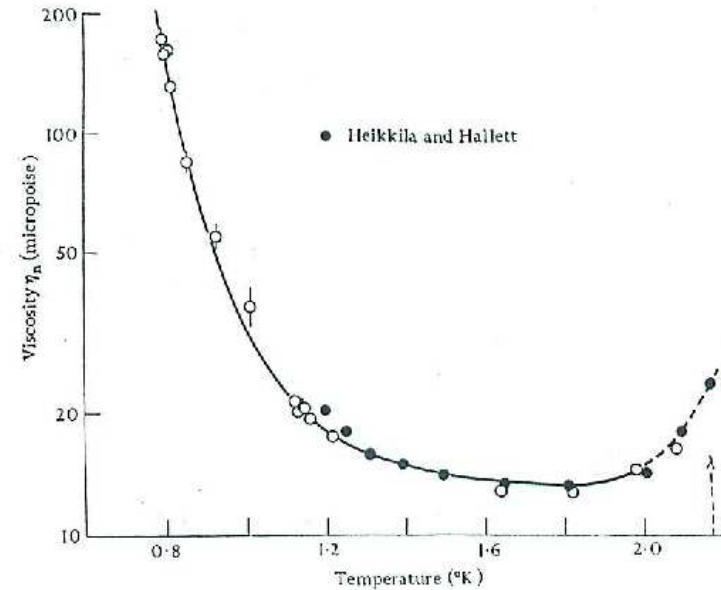


FIG. 11. The viscosity ( $\eta_n$ ) of helium II as measured in a rotation viscometer (Woods and Hollis Hallett [50]). The full points show the earlier results of Heikkila and Hollis Hallett [51].

Kapitza (1938)

viscosity vanishes below  $T_c$

capillary flow viscometer

Hollis-Hallett (1955)

roton minimum, phonon rise

rotation viscometer

$$\eta/s \simeq 0.8 \hbar/k_B$$

# Viscosity near Liquid-Gas Transition

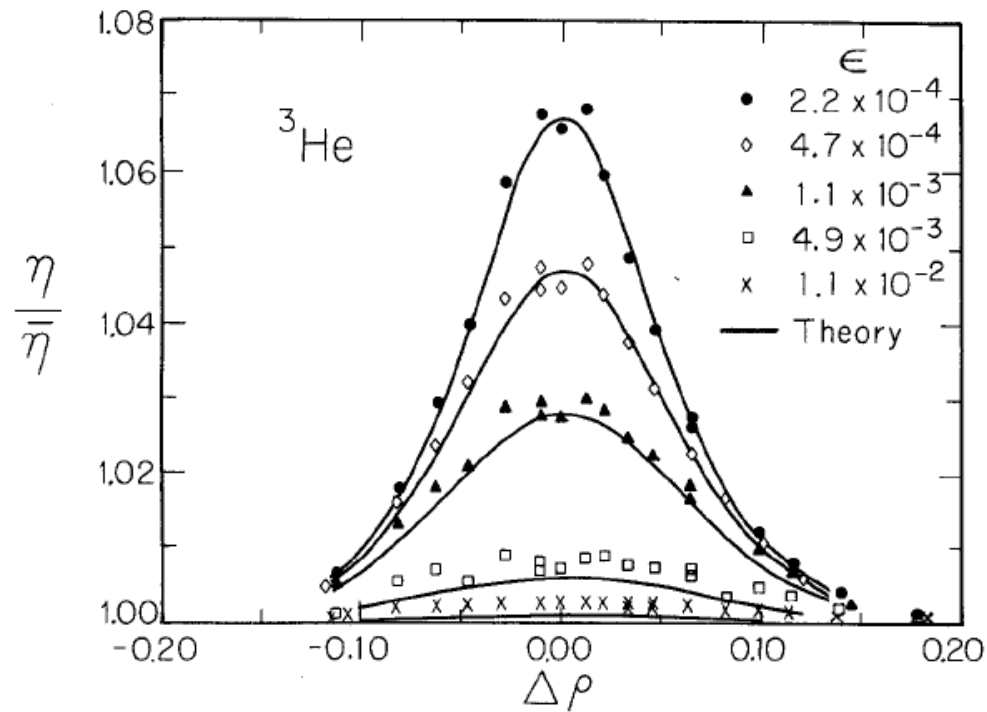
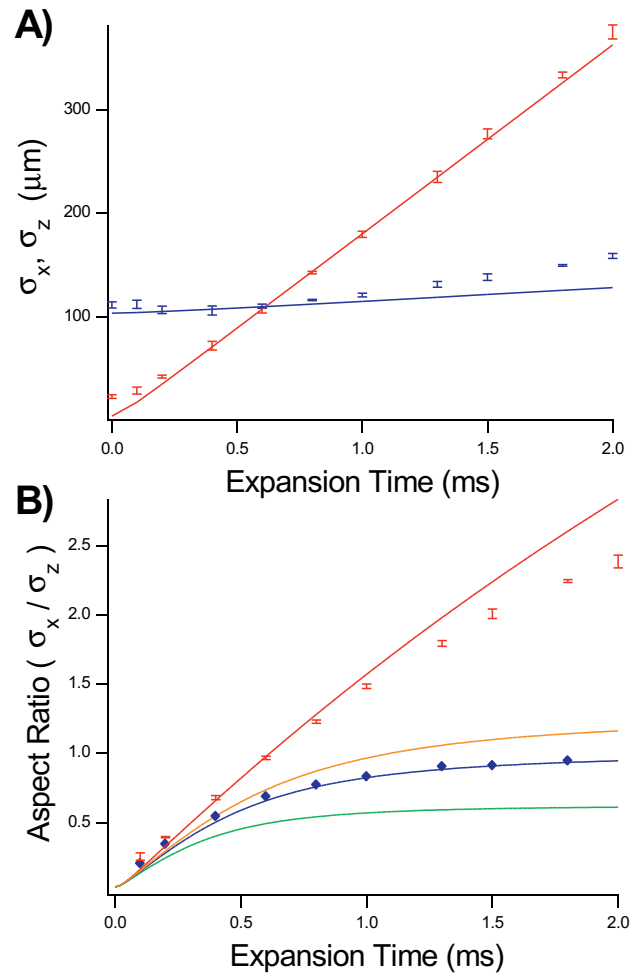
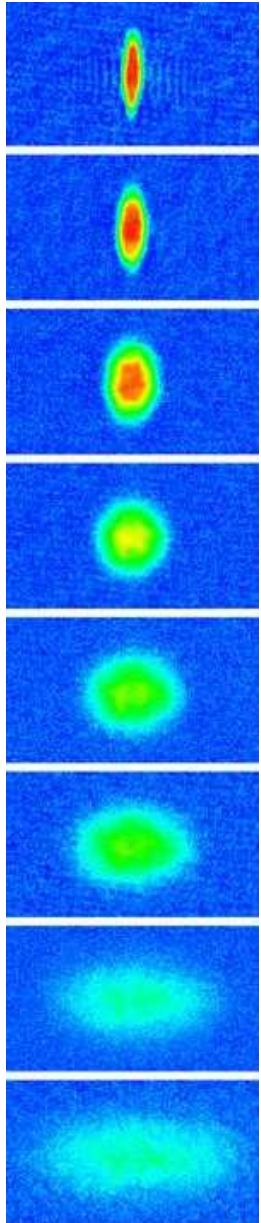


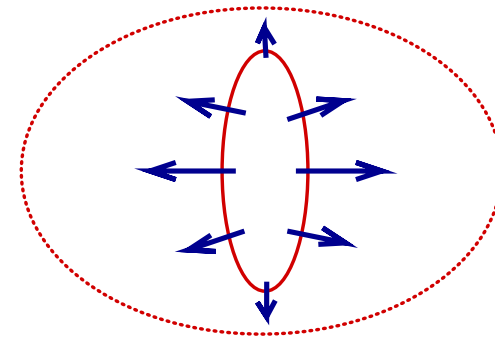
Fig. 14. The normalized viscosity  $\tilde{\eta}/\bar{\eta}$  for  ${}^3\text{He}$  versus  $\Delta\rho_c$  along several isotherms. (—) Fits of the MC theory with  $q_D = 3 \times 10^6 \text{ cm}^{-1}$ .

Agosta, Wang, Cohen, Meyer (1987)

## II. Elliptic Flow (Fermions)

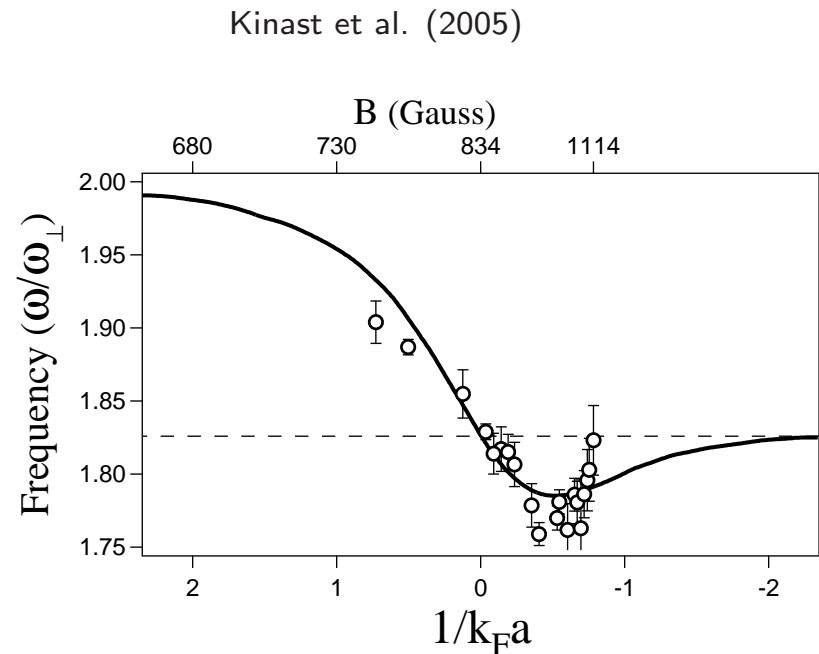
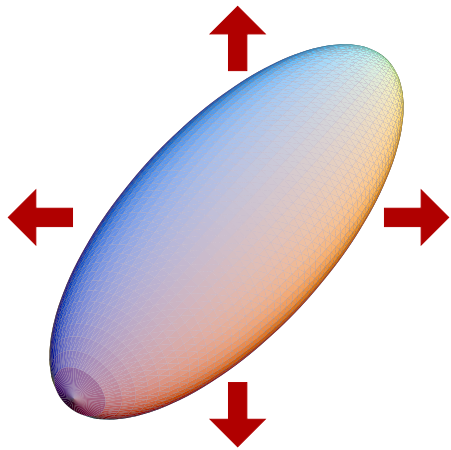


Hydrodynamic  
expansion converts  
coordinate space  
anisotropy  
to momentum space  
anisotropy



# Collective Modes

Radial breathing mode



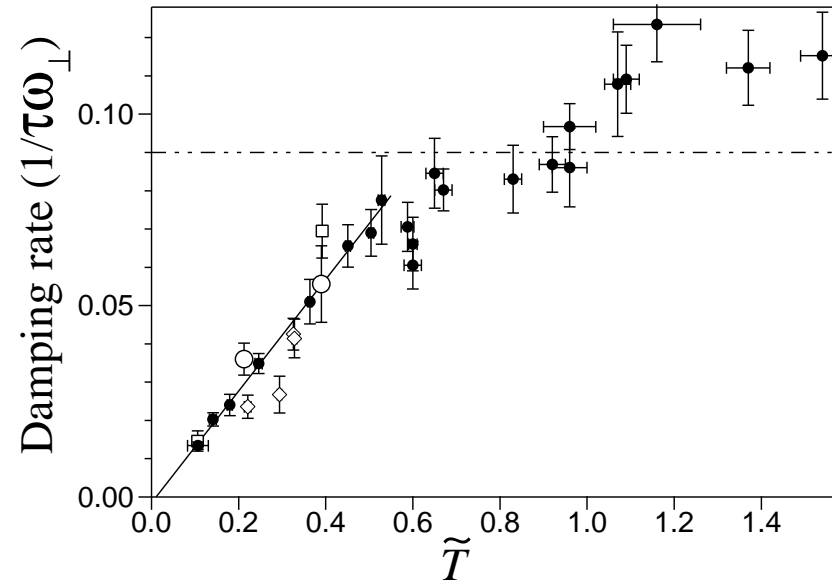
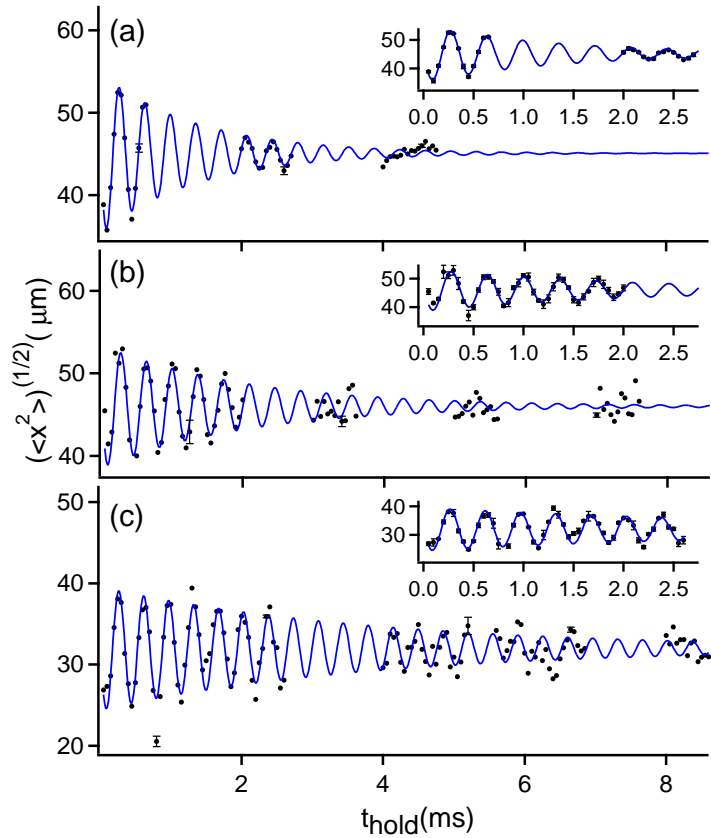
Ideal fluid hydrodynamics, equation of state  $P \sim n^{5/3}$

$$\frac{\partial n}{\partial t} + \vec{\nabla} \cdot (n\vec{v}) = 0$$

$$\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \vec{\nabla}) \vec{v} = -\frac{1}{mn} \vec{\nabla} P - \frac{1}{m} \vec{\nabla} V$$

$$\omega = \sqrt{\frac{10}{3}} \omega_{\perp}$$

# Damping of Collective Excitations



$$T/T_F = (0.5, 0.33, 0.17)$$

$\tau\omega$ : decay time  $\times$  trap frequency

Kinast et al. (2005)

# Viscous Hydrodynamics

Energy dissipation ( $\eta, \zeta, \kappa$ : shear, bulk viscosity, heat conductivity)

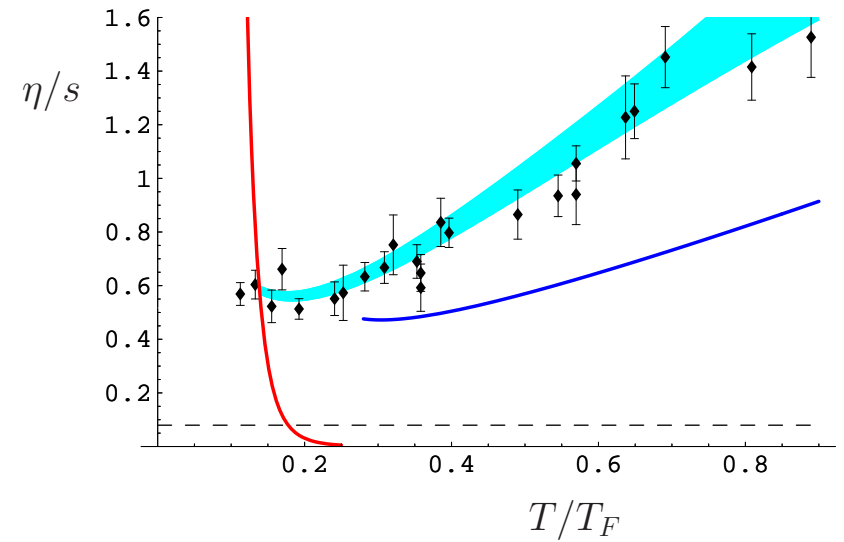
$$\dot{E} = -\frac{\eta}{2} \int d^3x \left( \partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \partial_k v_k \right)^2$$

$$- \zeta \int d^3x (\partial_i v_i)^2 - \frac{\kappa}{T} \int d^3x (\partial_i T)^2$$

Shear viscosity to entropy ratio  
(assuming  $\zeta = \kappa = 0$ )

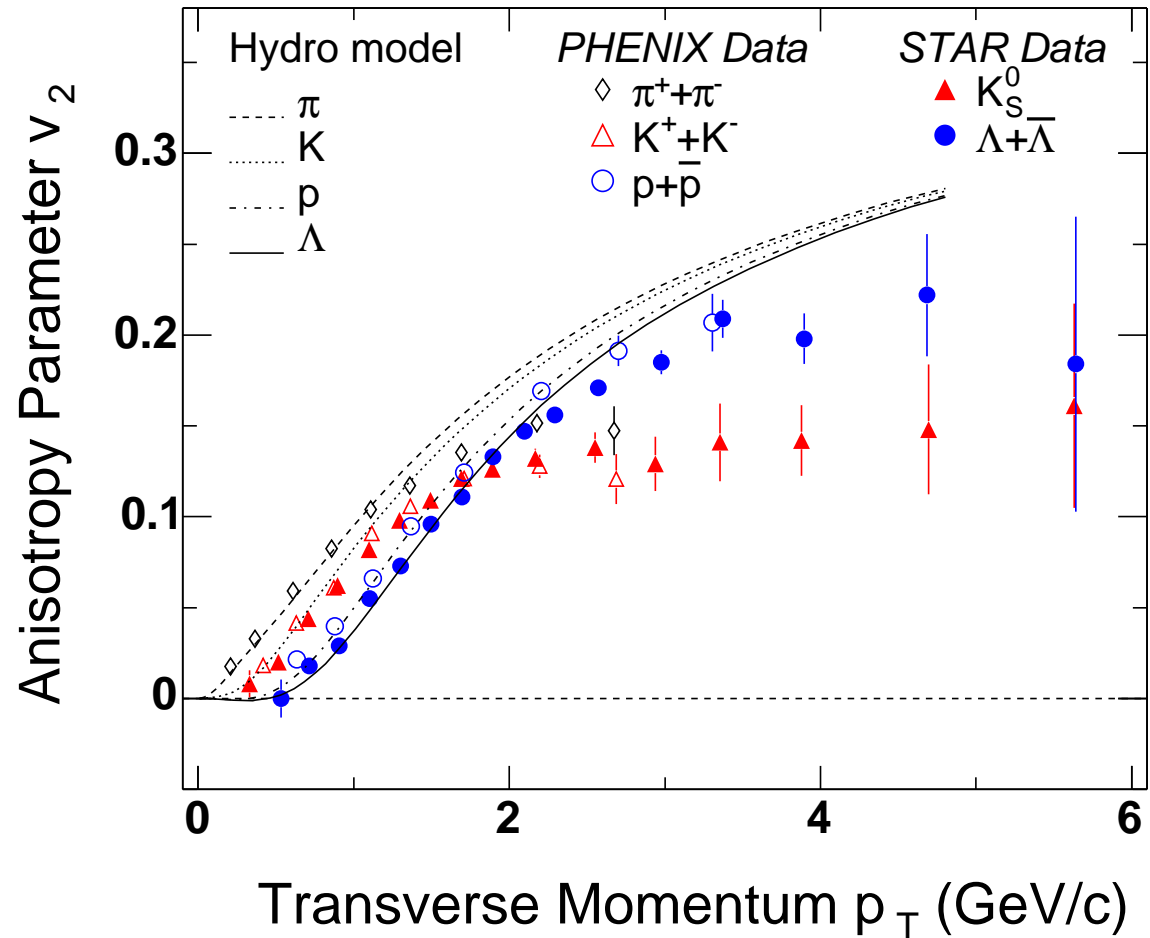
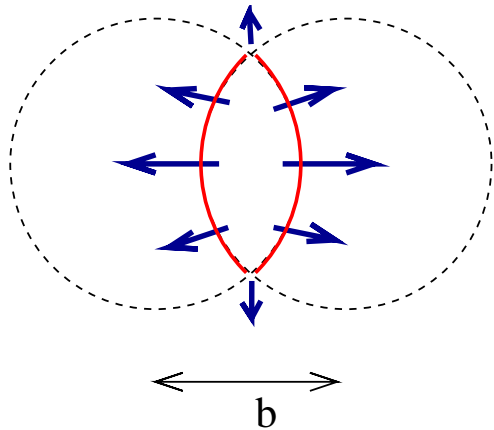
$$\frac{\eta}{s} = \frac{3}{4} \xi^{\frac{1}{2}} (3N)^{\frac{1}{3}} \frac{\Gamma}{\omega_{\perp}} \frac{\bar{\omega}}{\omega_{\perp}} \frac{N}{S}$$

see also Bruun, Smith, Gelman et al.



# III. Elliptic Flow (QGP)

Hydrodynamic expansion converts  
coordinate space  
anisotropy  
to momentum space  
anisotropy



source: U. Heinz (2005)

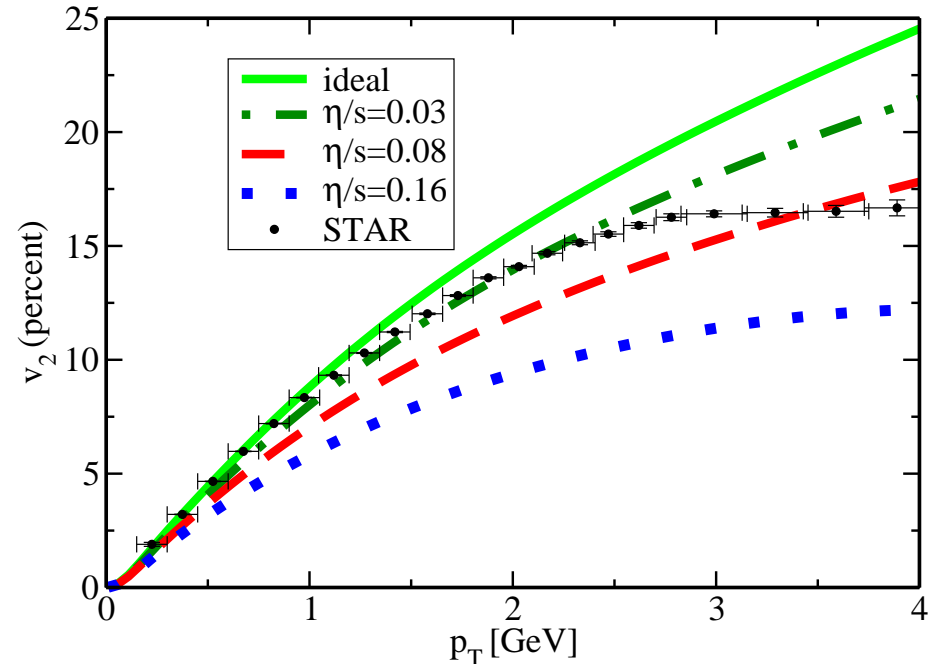
# Viscosity and Elliptic Flow

Consistency condition  $T_{\mu\nu} \gg \delta T_{\mu\nu}$   
(applicability of Navier-Stokes)

$$\frac{\eta + \frac{4}{3}\zeta}{s} \ll \frac{3}{4}(\tau T)$$

Danielewicz, Gyulassy (1985)

Very restrictive for  $\tau < 1$  fm



Romatschke (2007), Teaney (2003)

Many questions: Dependence on initial conditions, freeze out, etc.



## Outlook

Too early to declare a winner.

$$\eta/s \simeq 0.8 \text{ (He)}, \quad \eta/s \leq 0.5 \text{ (CA)}, \quad \eta/s \leq 0.5 \text{ (QGP)}$$

Other experimental constraints, more analysis needed.

Existing theory tools: Kinetic theory in low/high  $T$  limit, Kubo+field theory tools (lattice, etc), dynamic universality, ...

New theory tools: AdS/Cold Atom correspondence? Field theory approaches in cross over regime (large  $N$ , epsilon expansions, ...)